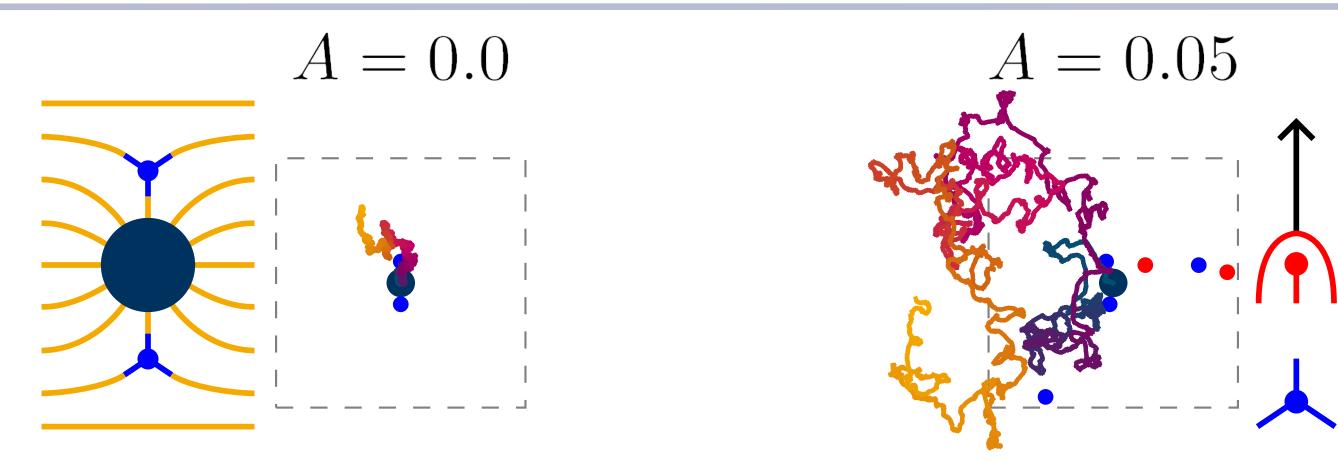


# DEFECT INTERACTIONS ENHANCE COLLOIDAL MOTILITY IN ACTIVE TURBULENCE

**Timofey Kozhukhov<sup>1,2</sup>** (timofey.kozhukhov@nbi.ku.dk), Kristian Thijssen<sup>2</sup>, Benjamin Loewe<sup>3</sup>, Tyler N. Shendruk<sup>1</sup>

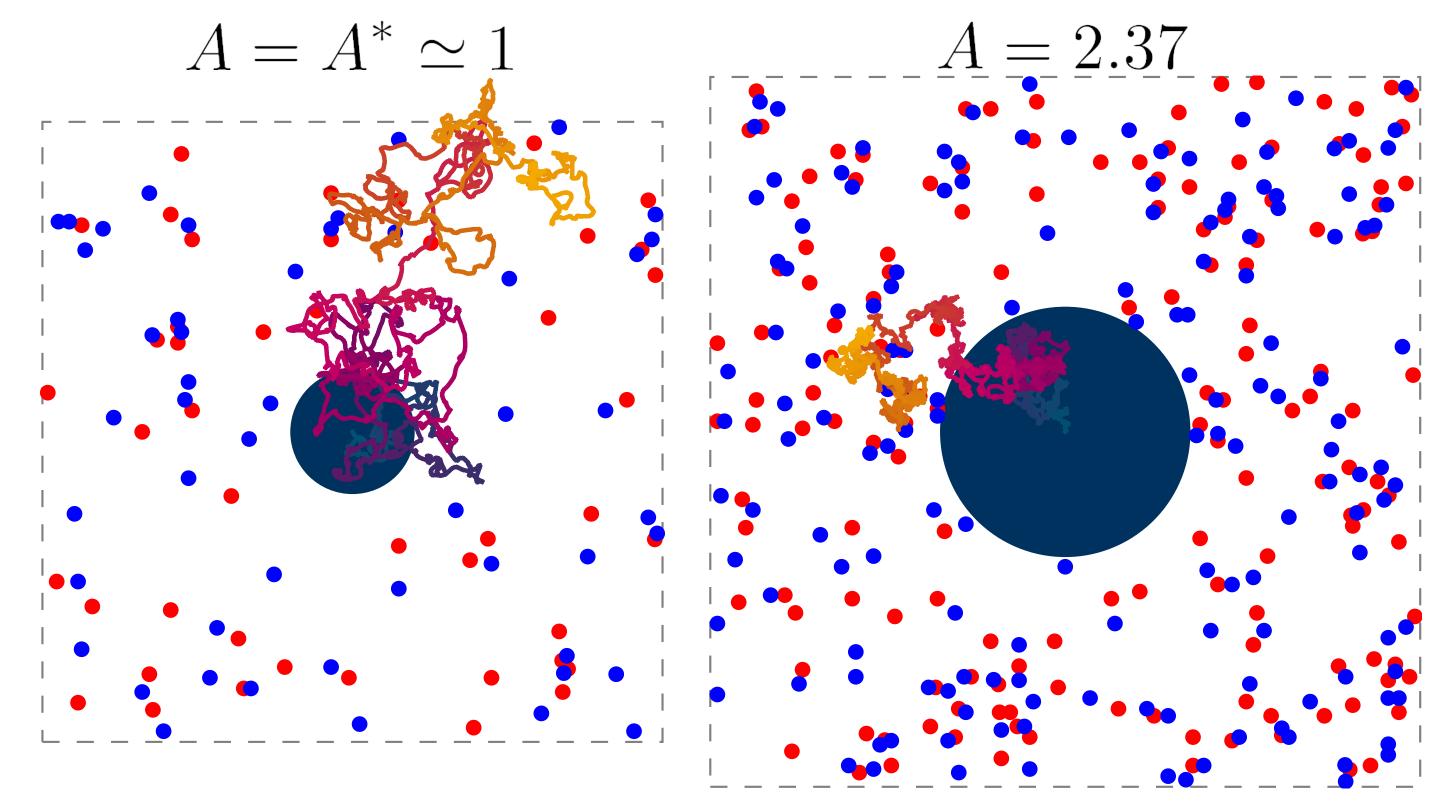
<sup>1</sup>School of Physics and Astronomy, The University of Edinburgh <sup>2</sup>Niels Bohr Institute, University of Copenhagen <sup>3</sup>Facultad de Física, Pontificia Universidad Católica de Chile





# Motivation

- There has been much recent interest on the behaviour of solutes in active-nematic turbulence [1, 2], which spontaneously produce pairs of advected -1/2 and selfmotile +1/2 defects.
- In passive systems (A = 0), strong nematic surface anchoring induces companion topological defects [3].
- The interplay between strongly anchored colloidal particles and active nematic tur-



### **Radial symmetry breaking**

# Solvent and solute model

- A two-dimensional extensile active-nematic solvent is simulated using AN-MPCD [4], a technique which replicates fluctuating active nematic turbulence.
- Colloidal disks are modelled as rigid boundary conditions with strong nematic surface anchoring that are free to move within the solvent [5].
- The effect of varying the colloidal radius  $R_{\rm C}$  and the active length scale  $\ell_{\alpha}$  is quantified through a dimensionless activity number A, defined as

 $A = \frac{\pi R_{\rm C}^2}{\pi \left(\ell_{\alpha}^2 + 2R_{\rm C}\ell_{\alpha}\right)} \sim \frac{R_{\rm C}^2}{\ell_{\alpha}^2}.$ 

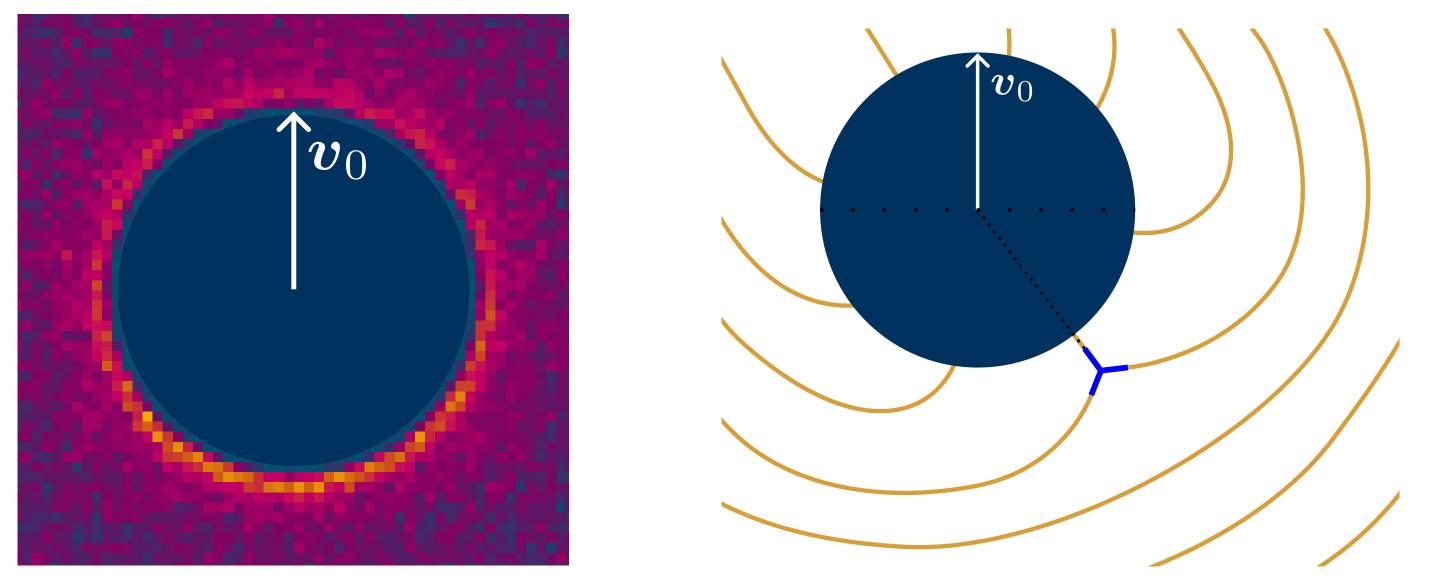
# Activity enhanced diffusion

• Differences in colloidal motility can be quantified through the mean squared displacement (MSD)  $\langle \Delta r^2 \rangle$  of colloidal trajectories r(t). They can be compared through a common MSD model of self-propelled particles

 $\langle \Delta r^2 \rangle (\delta t) = 2dD_{\text{eff}} \delta t - 2v_0^2 \tau_{\text{r}}^2 (1 - e^{-\delta t/\tau_{\text{r}}}); \ D_{\text{eff}} = D + v_0^2 \tau_{\text{r}}/d,$ 

- where d is the dimensionality,  $D_{\rm eff}$  is an effective diffusion coefficient that includes passive Brownian diffusion D, the instantaneous self-propulsion speed  $v_0$ , and the reorientation time  $\tau_{\rm r}$ .
- The instantaneous speed of the colloid  $v_0$  increases with activity A, platauing for

- This single remaining companion defect is found to lie behind the colloid, relative to the direction of motion.
- This leaves a colloid-companion complex with an effective +1/2 charge, which suggests it should be self-motile.
- However, the director pattern of this complex is in the opposite direction to the colloid's motion, inconsistent with +1/2 defects.

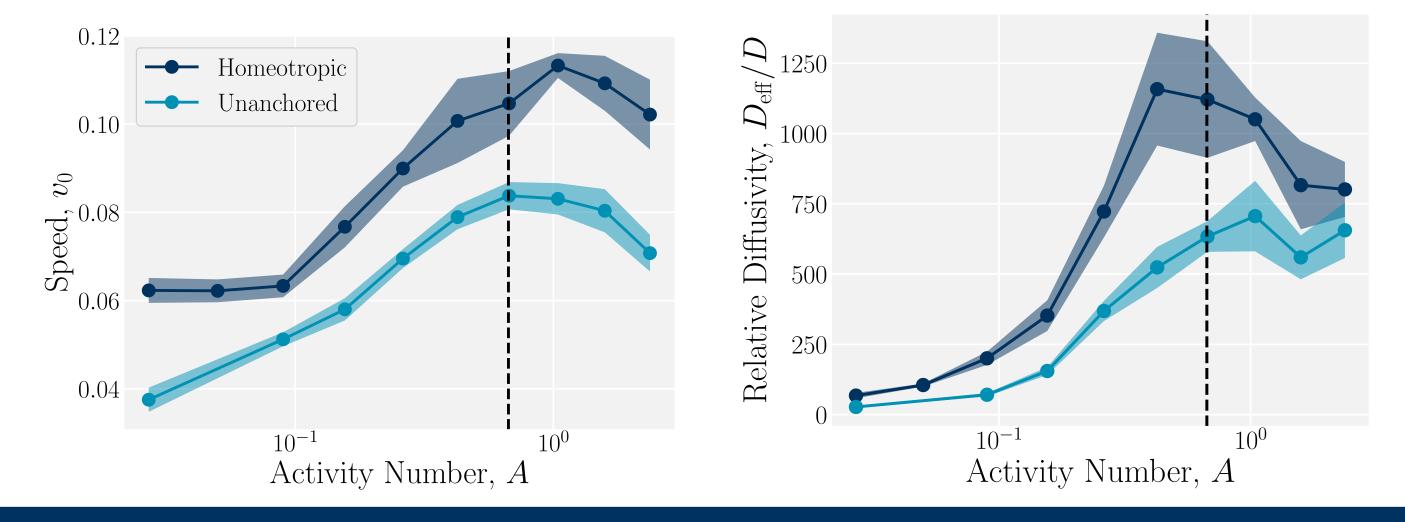


## **Mechanism for enhanced diffusion**

• Rather than the net charge of the colloid-companion complex, the enhanced motil-

 $A = A^* \simeq 0.67$ . Likewise, the relative diffusion  $D_{\rm eff}/D = 1 + D_{\alpha}/D$  observes a similar trend, but with a local maxima at  $A = A^*$ .

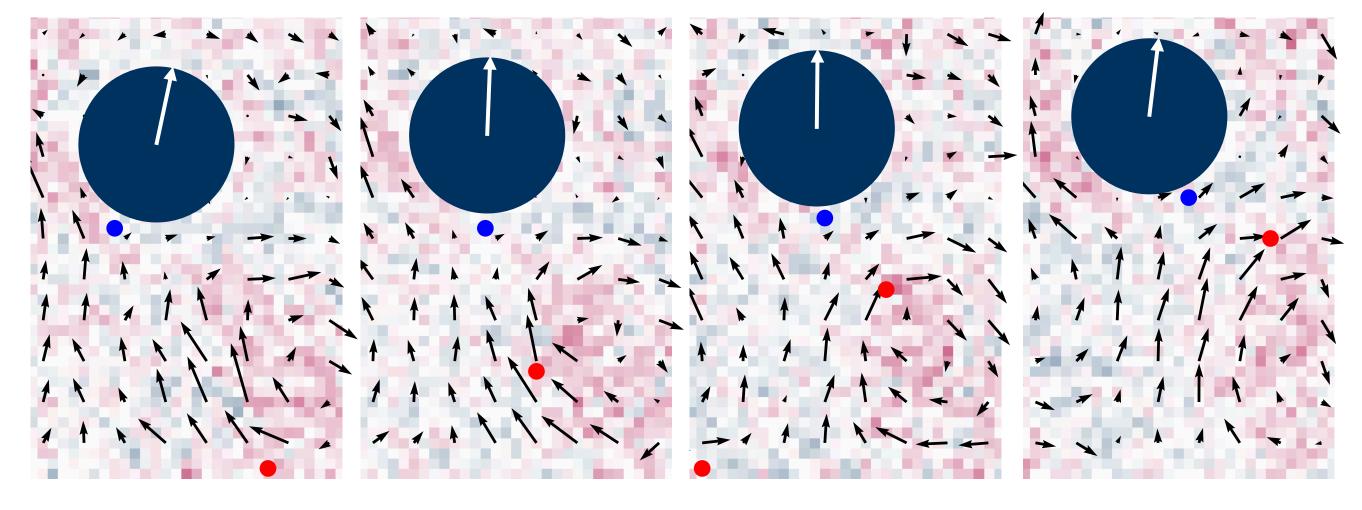
• However, unanchored colloids exhibit smaller speeds, and do not exhibit local maxima in the effective diffusion, which is instead monotonic.



### Changes in topological surroundings

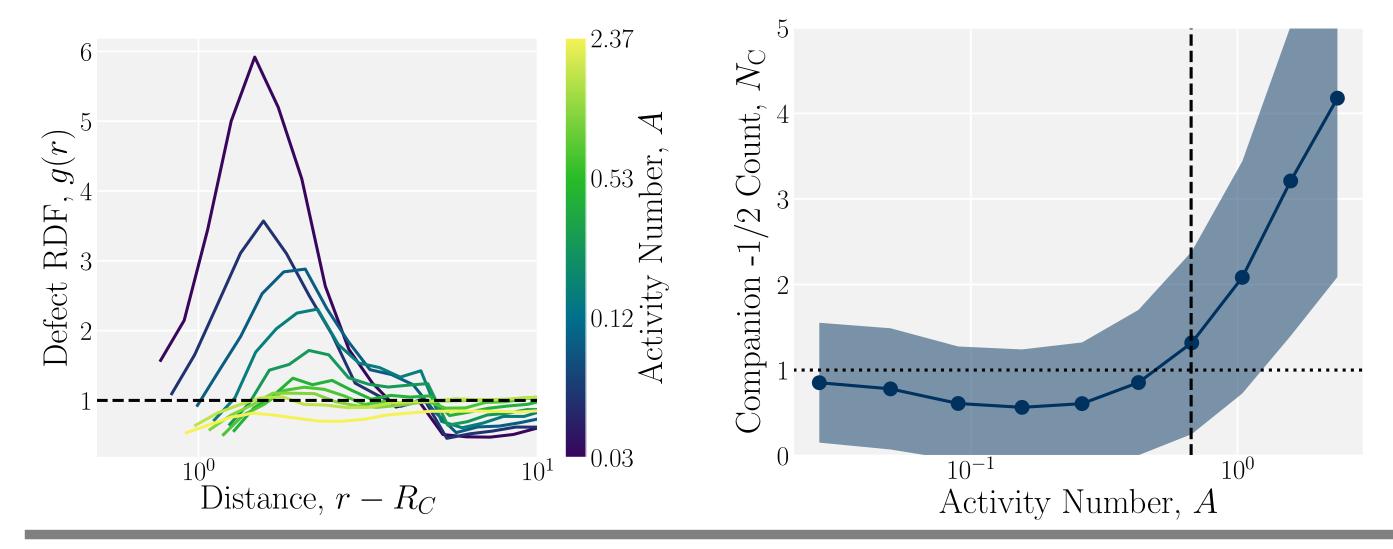
- The radial distribution function (RDF) of -1/2 defects describes the distribution of defects around the colloid.
- At low activities, the RDF exhibits a peak at short distances, indicative of companion defects. As activity rises towards  $A^*$ , companion defects become indistinguishable, becoming more homogeneously distributed akin to a gas.
- Unlike the passive case, where there are 2 companion defects, low activities exhibit 1 companion on average. For  $A \gtrsim A^*$ , the number of companions rises.

- ity of the colloid is attributed to the interplay between the complex and surrounding +1/2 defects.
- $\bullet + 1/2$  defects are attracted to the remaining companion defect, but are deflected due to the elastic forces from colloidal anchoring. The resulting flow fields push the colloid-companion complex away, resulting in enhanced motility.



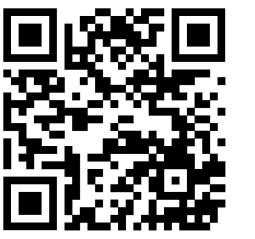
### **Conclusions and outlook**

- We have revealed a novel mechanism for the enhanced motility of strongly anchored colloids in active turbulence, dependent on a dimensionless activity number that compares the colloidal radius to the active length scale.
- For low activities  $A \leq A^*$ , colloids exhibit enhanced motility and non-monotonic diffusion due to +1/2 defects deflected by elastic forces.
- At higher activities  $A \gtrsim A^*$ , the homogeneous topological surroundings of the col-



loid balances any net flow from incoming +1/2 defects, saturating ballistic speed and decreasing effective diffusion.

• This work begins to extend the rich dynamics of anchored colloidal suspensions in passive nematics to active turbulence, opening the door to using colloids as a means of manipulating and controlling defect behaviour in active systems.







#### References

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View movies at www.kozhukhov.co.uk